

# HW1 Q1 Two-Path Channels

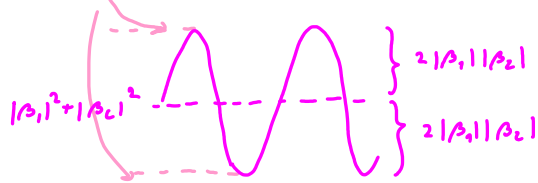
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In class, we have seen that

$$|H(f)|^2 = |\beta_1|^2 + |\beta_2|^2 + 2|\beta_1||\beta_2| \cos(2\pi(\Delta t_2 - \Delta t_1)f + (\angle\beta_1 - \angle\beta_2))$$

$$\max = |\beta_1|^2 + |\beta_2|^2 + 2|\beta_1||\beta_2|$$

$$\min = |\beta_1|^2 + |\beta_2|^2 - 2|\beta_1||\beta_2|$$

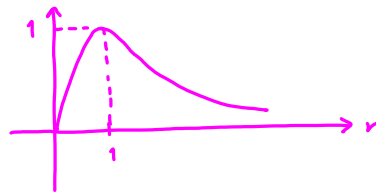


In frequency domain, the "frequency" of the oscillation is determined by  $\Delta t_2 - \Delta t_1$ .

Large  $\Delta t_2 - \Delta t_1 \rightarrow$  more oscillation  $\rightarrow$  (i), (iii)

Small  $\Delta t_2 - \Delta t_1 \rightarrow$  less oscillation  $\rightarrow$  (ii), (iv)

$$\text{Depth of fluctuation} = \frac{4|\beta_1||\beta_2|}{|\beta_1|^2 + |\beta_2|^2 + 2|\beta_1||\beta_2|} = \frac{4}{\frac{|\beta_1|}{|\beta_2|} + \frac{|\beta_2|}{|\beta_1|} + 2} = \frac{4}{r + \frac{1}{r} + 2}$$



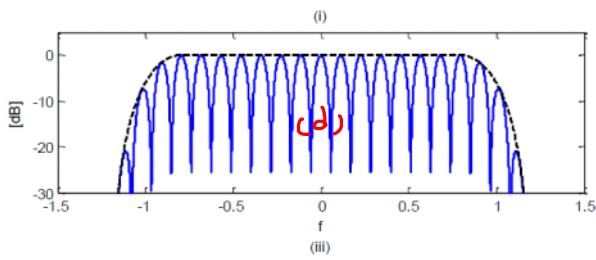
$$r = \frac{|\beta_1|}{|\beta_2|}$$

So, depth is large when  $|\beta_1| \approx |\beta_2|$

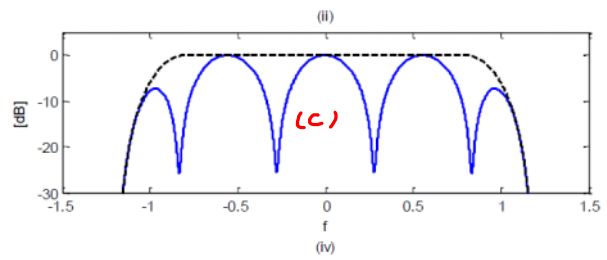
$\rightarrow$  (i), (ii)

small when  $|\beta_1| \gg |\beta_2|$  or  $|\beta_1| \ll |\beta_2| \rightarrow$  (iii), (iv)

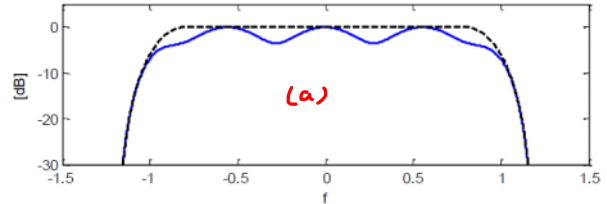
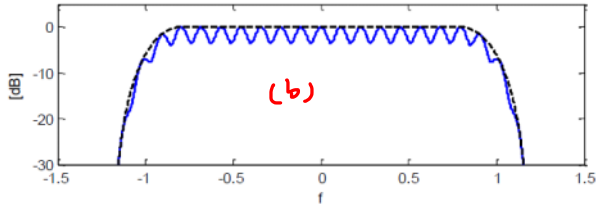
Large  $\Delta t_2 - \Delta t_1$



small  $\Delta t_2 - \Delta t_1$



$|\beta_1| \approx |\beta_2|$



$|\beta_1| \gg |\beta_2|$   
or  
 $|\beta_2| \gg |\beta_1|$

# HW1 Q2 GSOP for Complex-Valued Vectors

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From MATLAB

(a)  $\vec{e}^{(1)} = \frac{1}{2} \begin{pmatrix} 1+j \\ 1-j \\ 0 \end{pmatrix}, \vec{e}^{(2)} = \frac{1}{2} \begin{pmatrix} 1+j \\ -1+j \\ 0 \end{pmatrix}, \vec{e}^{(3)} = \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix}$

(b)  $C = \begin{bmatrix} 2 & -j & 1 & -1 \\ 0 & 1 & -j & j \\ 0 & 0 & 1 & j \end{bmatrix}$

## Detailed Solution

(a)  $\vec{u}^{(1)} = \vec{v}^{(1)} = \begin{pmatrix} 1+j \\ 1-j \\ 0 \end{pmatrix}, \|\vec{u}^{(1)}\|^2 = (1^2+1^2) + (1^2+(-1)^2) + 0 = 2+2 = 4$   
 $\|\vec{u}^{(1)}\| = 2$

$\vec{e}^{(1)} = \frac{\vec{u}^{(1)}}{\|\vec{u}^{(1)}\|} = \frac{1}{2} \begin{pmatrix} 1+j \\ 1-j \\ 0 \end{pmatrix}$

$\vec{u}^{(2)} = \vec{v}^{(2)} - \langle \vec{v}^{(2)}, \vec{e}^{(1)} \rangle \vec{e}^{(1)} = \begin{pmatrix} 1 \\ -1 \\ 0 \end{pmatrix} - (-j) \frac{1}{2} \begin{pmatrix} 1+j \\ 1-j \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ -1 \\ 0 \end{pmatrix} + \frac{1}{2} \begin{pmatrix} j-1 \\ j+1 \\ 0 \end{pmatrix}$   
 $\left(\frac{1}{2}\right)^* \left( (1-j) - (1+j) \right)$  Note that the  $j$  is conjugated.  $= \frac{1}{2} \begin{pmatrix} j+1 \\ j-1 \\ 0 \end{pmatrix}$   
 $= \frac{1}{2} (-2j) = -j$

$\|\vec{u}^{(2)}\|^2 = \frac{1}{2} \left( (1^2+1^2) + (1^2+(-1)^2) + 0 \right)$

$\vec{e}^{(2)} = \frac{\vec{u}^{(2)}}{\|\vec{u}^{(2)}\|} = \vec{u}^{(2)} = \frac{1}{2} \begin{pmatrix} j+1 \\ j-1 \\ 0 \end{pmatrix} = \frac{1}{2} (4) = 1$

$\vec{u}^{(3)} = \vec{v}^{(3)} - \langle \vec{v}^{(3)}, \vec{e}^{(1)} \rangle \vec{e}^{(1)} - \langle \vec{v}^{(3)}, \vec{e}^{(2)} \rangle \vec{e}^{(2)}$   
 $\left(\frac{1}{2}\right)^* \left( (1-j) + (1+j) + 0 \right) = \frac{1}{2} (2) = 1$   
 $\left(\frac{1}{2}\right)^* \left( -j+1-j-1 \right) = \frac{1}{2} (-2j) = -j$

$= \begin{pmatrix} 1 \\ 1 \\ -1 \end{pmatrix} - \frac{1}{2} \begin{pmatrix} 1+j \\ 1-j \\ 0 \end{pmatrix} + j \frac{1}{2} \begin{pmatrix} j+1 \\ j-1 \\ 0 \end{pmatrix} = \frac{1}{2} \left( \begin{pmatrix} 2 & -1-j & + & -1+j \\ 2 & -1+j & + & -1-j \\ -2 & +0 & + & 0 \end{pmatrix} \right) = \frac{1}{2} \begin{pmatrix} 0 \\ 0 \\ -2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix}$

$\|\vec{u}^{(3)}\| = \sqrt{0^2+0^2+(-1)^2} = 1$

$\vec{e}^{(3)} = \frac{\vec{u}^{(3)}}{\|\vec{u}^{(3)}\|} = \vec{u}^{(3)} = \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix}$

$\vec{u}^{(4)} = \vec{v}^{(4)} - \langle \vec{v}^{(4)}, \vec{e}^{(1)} \rangle \vec{e}^{(1)} - \langle \vec{v}^{(4)}, \vec{e}^{(2)} \rangle \vec{e}^{(2)} - \langle \vec{v}^{(4)}, \vec{e}^{(3)} \rangle \vec{e}^{(3)}$

$$\begin{aligned}
 u &= v - \underbrace{\langle v, e \rangle e}_{\substack{= \left(\frac{1}{2}\right)^* (-(-1-j) - (1+j)) \\ = \frac{1}{2} (-1+j-1-j) \\ = \frac{-2}{2} = -1}} - \underbrace{\langle v, e \rangle e}_{\substack{= \left(\frac{1}{2}\right)^* (-(j+1) - (-j-1)) \\ = \frac{1}{2} (j-1+j+1) \\ = \frac{1}{2} (2j) = j}} - \underbrace{\langle v, e \rangle e}_{= (-j)(-1) = j} \\
 &= \frac{1}{2} \begin{pmatrix} -2 & +1+j & -(-1+j) & +0 \\ -2 & +1-j & -(-1-j) & +0 \\ -2j & +0 & -0 & +2j \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \leftarrow \text{so, no } \vec{u}^{(4)} \text{ and hence no } \vec{e}^{(4)}.
 \end{aligned}$$

(b) From the procedure above, we get

$$\vec{u}^{(1)} = \vec{v}^{(1)} \Rightarrow \vec{v}^{(1)} = \vec{u}^{(1)} = \|\vec{u}^{(1)}\| \vec{e}^{(1)} = 2 \vec{e}^{(1)} = E \begin{pmatrix} 2 \\ 0 \\ 0 \end{pmatrix}$$

$$\vec{u}^{(2)} = \vec{v}^{(2)} - (-j) \vec{e}^{(1)} \Rightarrow \vec{v}^{(2)} = (-j) \vec{e}^{(1)} + \|\vec{u}^{(2)}\| \vec{e}^{(2)} = (-j) \vec{e}^{(1)} + 1 \vec{e}^{(2)} = E \begin{pmatrix} -j \\ 1 \\ 0 \end{pmatrix}$$

$$\begin{aligned}
 \vec{u}^{(3)} &= \vec{v}^{(3)} - 1 \vec{e}^{(1)} - (-j) \vec{e}^{(2)} \Rightarrow \vec{v}^{(3)} = (1) \vec{e}^{(1)} + (-j) \vec{e}^{(2)} + \|\vec{u}^{(3)}\| \vec{e}^{(3)} \\
 &= (1) \vec{e}^{(1)} + (-j) \vec{e}^{(2)} + (1) \vec{e}^{(3)} = E \begin{pmatrix} 1 \\ -j \\ 1 \end{pmatrix}
 \end{aligned}$$

$$\vec{u}^{(4)} = \vec{v}^{(4)} - (-1) \vec{e}^{(1)} - (j) \vec{e}^{(2)} - (j) \vec{e}^{(3)} \Rightarrow \vec{v}^{(4)} = (-1) \vec{e}^{(1)} + (j) \vec{e}^{(2)} + (j) \vec{e}^{(3)} = E \begin{pmatrix} -1 \\ j \\ j \end{pmatrix}$$

$$\text{so, } V = [\vec{v}^{(1)} \quad \vec{v}^{(2)} \quad \vec{v}^{(3)} \quad \vec{v}^{(4)}] = E \underbrace{\begin{bmatrix} 2 & -j & 1 & -1 \\ 0 & 1 & -j & j \\ 0 & 0 & 1 & j \end{bmatrix}}_C.$$

# HW1 Q3 Signal Space and Constellation

Monday, July 08, 2013 10:39 AM

(a)

$$\varepsilon_1 = \varepsilon_{s_1} = \int_{-\infty}^{\infty} |s_1(t)|^2 dt = \int_0^{T_b} v^2 dt = v^2 T_b \equiv \varepsilon$$

$$\varepsilon_2 = \varepsilon_{s_2} = \int_{-\infty}^{\infty} |s_2(t)|^2 dt = \int_0^{T_b} v^2 dt = v^2 T_b \equiv \varepsilon$$

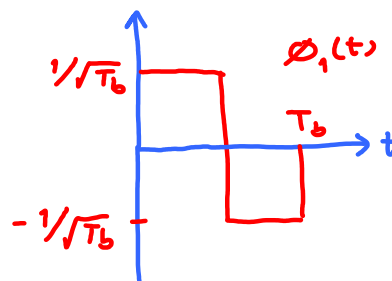
(b)

$$u_1(t) = s_1(t)$$

$$\phi_1(t) = \frac{s_1(t)}{\sqrt{\varepsilon_1}} = \frac{1}{\sqrt{\varepsilon}} s_1(t)$$

$$\frac{v}{\sqrt{\varepsilon}} = \frac{v}{v\sqrt{T_b}} = \frac{1}{\sqrt{T_b}}$$

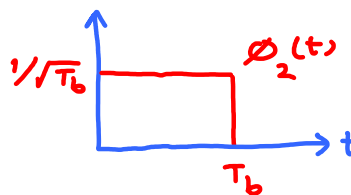
$$= \begin{cases} 1/\sqrt{T_b}, & 0 \leq t \leq \frac{T_b}{2} \\ -1/\sqrt{T_b}, & \frac{T_b}{2} < t < T_b \\ 0, & \text{otherwise} \end{cases}$$



$$u_2(t) = s_2(t) - \text{proj}_{u_1} s_2 = s_2(t) - \frac{\langle s_2, s_1 \rangle}{\langle s_1, s_1 \rangle} s_1(t) = s_2(t)$$

$$\phi_2(t) = \frac{u_2(t)}{\sqrt{\varepsilon_{u_2}}} = \frac{s_2(t)}{\sqrt{\varepsilon_2}} = \frac{1}{\sqrt{\varepsilon}} s_2(t)$$

$u_2 = s_2$

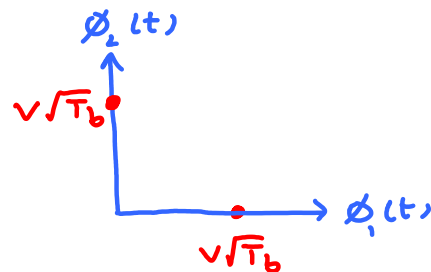


$$= \begin{cases} 1/\sqrt{T_b}, & 0 \leq t \leq T_b \\ 0, & \text{otherwise} \end{cases}$$

(c)

$$s_1(t) = \sqrt{\varepsilon} \phi_1(t) \Rightarrow \vec{s}^{(1)} = \begin{pmatrix} \sqrt{\varepsilon} \\ 0 \end{pmatrix} = \begin{pmatrix} v\sqrt{T_b} \\ 0 \end{pmatrix}$$

$$s_2(t) = \sqrt{\varepsilon} \phi_2(t) \Rightarrow \vec{s}^{(2)} = \begin{pmatrix} 0 \\ \sqrt{\varepsilon} \end{pmatrix} = \begin{pmatrix} 0 \\ v\sqrt{T_b} \end{pmatrix}$$



# HW1 Q4 Signal Space and Constellation

Monday, July 08, 2013 11:17 AM

(a)

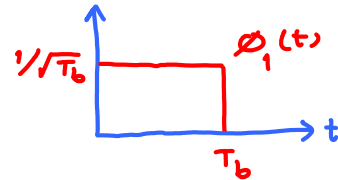
$$\begin{aligned} \mathcal{E}_1 = \mathcal{E}_{s_1} &= \int_{-\infty}^{\infty} |s_1(t)|^2 dt = \int_0^{T_b} v^2 dt = v^2 T_b. \\ \mathcal{E}_2 = \mathcal{E}_{s_2} &= \int_{-\infty}^{\infty} |s_2(t)|^2 dt = \int_0^{T_b} v^2 dt = v^2 T_b. \end{aligned}$$

↙ ≡  $\mathcal{E}$   
↘ ≡  $\mathcal{E}$

(b)

$$u_1(t) = s_1(t).$$

$$\phi_1(t) = \frac{s_1(t)}{\sqrt{\mathcal{E}_1}} = \frac{1}{\sqrt{\mathcal{E}}} s_1(t) = \begin{cases} 1/\sqrt{T_b}, & 0 < t < T_b \\ 0, & \text{otherwise.} \end{cases}$$



$$\frac{v}{\sqrt{\mathcal{E}}} = \frac{v}{v\sqrt{T_b}} = \frac{1}{\sqrt{T_b}}$$

$$u_2(t) = s_2(t) - \text{proj}_{u_1} s_2 = s_2(t) - \frac{\langle s_2, s_1 \rangle}{\langle s_1, s_1 \rangle} s_1(t)$$

$$\langle s_2, s_1 \rangle = \int_{-\infty}^{\infty} s_2(t) s_1(t) dt$$

$$= \alpha v^2 - (T_b - \alpha) v^2 = 2\alpha v^2 - T_b v^2$$

$$u_2(t) = s_2(t) - \frac{2\alpha v^2 - T_b v^2}{v^2 T_b} s_1(t) = s_2(t) - \left( \frac{2\alpha}{T_b} - 1 \right) s_1(t)$$

$$= s_2(t) + \left( 1 - \frac{2\alpha}{T_b} \right) s_1(t)$$

When  $\alpha = \frac{T_b}{2}$ , this term = 0 and  $u_2(t) = s_2(t)$ .

When  $\alpha < \frac{T_b}{2}$ , this term is  $> 0$ . So,  $s_2(t)$  will be shifted up.

When  $\alpha > \frac{T_b}{2}$ , this term is  $< 0$ . So,  $s_2(t)$  will be shifted down.

$$\int v + \left( 1 - \frac{2\alpha}{T_b} \right) v \quad \text{for } 0 < t < \alpha,$$

$$= \begin{cases} v + (1 - \frac{2\alpha}{T_b})v & \text{for } 0 < t < \alpha, \\ -v + (1 - \frac{2\alpha}{T_b})v & \text{for } \alpha < t < T_b, \\ 0 & \text{otherwise.} \end{cases}$$

$$= \begin{cases} (2 - \frac{2\alpha}{T_b})v & \text{for } 0 < t < \alpha, \\ -\frac{2\alpha}{T_b}v & \text{for } \alpha < t < T_b, \\ 0 & \text{otherwise.} \end{cases} = 2v \times \begin{cases} 1 - \frac{\alpha}{T_b} & \text{for } 0 < t < \alpha, \\ -\frac{\alpha}{T_b} & \text{for } \alpha < t < T_b, \\ 0 & \text{otherwise.} \end{cases}$$

Note: The shift amount is just enough to make  $\int u_2 = 0$

$$\left[ (1 - \frac{\alpha}{T_b})\alpha + (-\frac{\alpha}{T_b})(T_b - \alpha) = \alpha - \frac{\alpha^2}{T_b} - \alpha + \frac{\alpha^2}{T_b} = 0. \right]$$

$$\begin{aligned} \epsilon_{m_2} &= 4v^2 \left( \underbrace{\left(1 - \frac{\alpha}{T_b}\right)^2 \alpha + \left(-\frac{\alpha}{T_b}\right)^2 (T_b - \alpha)}_{1 - \frac{2\alpha}{T_b} + \frac{\alpha^2}{T_b^2}} \right) = 4v^2 \left( \alpha - \frac{2\alpha^2}{T_b} + \frac{\alpha^3}{T_b^2} + \frac{\alpha^2}{T_b} - \frac{\alpha^3}{T_b^2} \right) \\ &= 4v^2 \alpha \left(1 - \frac{\alpha}{T_b}\right) \end{aligned}$$

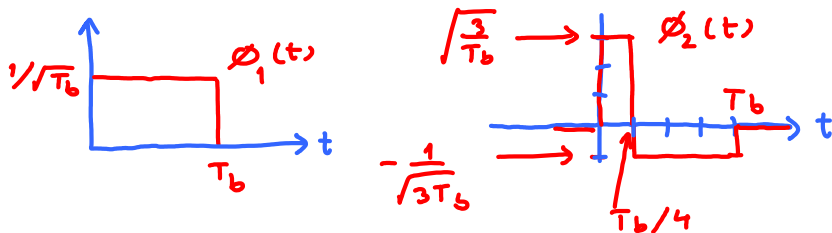
$$\Rightarrow \sqrt{\epsilon_{m_2}} = 2v \sqrt{\alpha \left(1 - \frac{\alpha}{T_b}\right)}$$

$$\phi_2(t) = \frac{u_2(t)}{\sqrt{\epsilon_{m_2}}} = \frac{2v}{2v \sqrt{\alpha \left(1 - \frac{\alpha}{T_b}\right)}} \begin{cases} 1 - \frac{\alpha}{T_b} & \text{for } 0 < t < \alpha, \\ -\frac{\alpha}{T_b} & \text{for } \alpha < t < T_b, \\ 0 & \text{otherwise.} \end{cases}$$

$$= \frac{1}{\sqrt{\alpha \left(1 - \frac{\alpha}{T_b}\right)}} \times \begin{cases} 1 - \frac{\alpha}{T_b} & \text{for } 0 < t < \alpha, \\ -\frac{\alpha}{T_b} & \text{for } \alpha < t < T_b, \\ 0 & \text{otherwise.} \end{cases}$$

(C) When  $\alpha = \frac{1}{4}T_b$ ,  $\Rightarrow \frac{1}{\sqrt{\frac{T_b}{4} \left(1 - \frac{1}{4}\right)}} = \frac{4}{\sqrt{3}T_b}$

$$\phi_2(t) = \frac{4}{\sqrt{3}T_b} \times \begin{cases} 3/4 & \text{for } 0 < t < \frac{T_b}{4} \\ -1/4 & \text{for } \frac{T_b}{4} < t < T_b \\ 0 & \text{otherwise} \end{cases} = \begin{cases} \sqrt{3}/T_b, & \text{for } 0 < t < \frac{T_b}{4}, \\ -\frac{1}{\sqrt{3}T_b}, & \text{for } \frac{T_b}{4} < t < T_b, \\ 0, & \text{otherwise.} \end{cases}$$



(d)

$$s_1(t) = \sqrt{\epsilon_1} \phi_1(t) = \sqrt{\epsilon} \phi_1(t) = v\sqrt{T_b} \phi_1(t) \Rightarrow \vec{s}^{(1)} = \begin{pmatrix} v\sqrt{T_b} \\ 0 \end{pmatrix} = \begin{pmatrix} \sqrt{\epsilon} \\ 0 \end{pmatrix}$$

$$s_2(t) = u_2(t) - \left(1 - 2\frac{\alpha}{T_b}\right) s_1 = \sqrt{\epsilon_{m_2}} \phi_2(t) - \left(1 - 2\frac{\alpha}{T_b}\right) \sqrt{\epsilon} \phi_1(t)$$

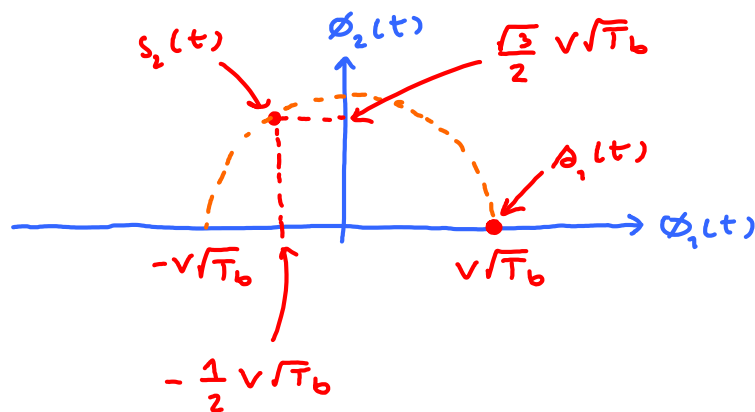
$$= 2v\sqrt{\alpha\left(1 - \frac{\alpha}{T_b}\right)} \phi_2(t) - \left(1 - 2\frac{\alpha}{T_b}\right) v\sqrt{T_b} \phi_1(t)$$

$$\Rightarrow \vec{s}^{(2)} = \begin{pmatrix} -\left(1 - 2\frac{\alpha}{T_b}\right) \\ 2\sqrt{\frac{\alpha}{T_b}\left(1 - \frac{\alpha}{T_b}\right)} \end{pmatrix} v\sqrt{T_b} = \sqrt{\epsilon} \begin{pmatrix} 2r-1 \\ 2\sqrt{r(1-r)} \end{pmatrix} \text{ where } r = \frac{\alpha}{T_b}$$

(e)

When  $\alpha = \frac{T_b}{4}$ ,  $s_2(t) = 2v\sqrt{\frac{T_b}{4}\left(\frac{3}{4}\right)} \phi_2(t) - \left(1 - \frac{1}{2}\right) v\sqrt{T_b} \phi_1(t)$

$$= \frac{\sqrt{3}}{2} v\sqrt{T_b} \phi_2(t) - \frac{1}{2} v\sqrt{T_b} \phi_1(t)$$



(f)

Note that  $(2r-1)^2 + (2\sqrt{r(1-r)})^2 = 4r^2 - 4r + 1 + 4r - 4r^2 = 1$ .

$$r = \frac{\alpha}{T_b} = \frac{k}{10}$$

